
The Effect of The Symmetrical Elastic Nonlinearity on Structural Vibrations Transmitted by Dynamic Equipment on The Construction Envelope

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Abstract: - The results of the research carried out for elastic systems with symmetrical nonlinearity have been systematized in this article so that they can be disseminated for real models of dynamic equipment in construction endowment. Thus, the ventilation - conditioning systems, water pumping systems, main and interior piping systems from nuclear power plants, platforms with optoelectronic equipment are supported on elastic devices that can have a non-linear behavior.

In this case, it is highlighted that under certain functional and constructive conditions it is possible the occurrence of subharmonic vibrations of order 1/3 of the excitation pulsation. This leads to the generation of structural noise with consequences for the people inside it.

Keywords: - nonlinear vibrations, subharmonic vibrations, symmetrical elasticity

1. INTRODUCTION

It is adopted the model of the system with a degree of freedom with viscoelastic connection where the elastic force is a nonlinear function of the third order in relation to the instantaneous deformation.

In this case, in order to highlight the occurrence of 1/3 subharmonics order, it was used the method of external excitation on the frequency band sought for so that the mechanical work of external excitation be equal to the mechanical work of dissipation. [1-5]

2. EVALUATION OF SUBHARMONIC VIBRATIONS OF ORDER 1/3 OF THE EXCITATION PULSATION

The system is characterized by a coefficient of rigidity expressed by the relationship:

$$k=k_0(1+\beta q^2) \quad (1)$$

where for $\beta > 0$, the relation (1) characterizes a "strong" elastic element, and for $\beta < 0$ - a "soft" elastic element. It is mentioned that the parameter β is measured in m^{-2} , and k_0 in Nm^{-1} . The differential equation of the system is:

$$a\ddot{q} + b\dot{q} + k_0(1 + \beta q^2)q = Q_0 \cos \omega t \quad (2)$$

where q is the generalized coordinate which can be linear displacement or angular displacement; a - coefficient of inertia (mass m or moment of inertia J); b - viscous amortization factor; k_0 - the rigidity coefficient of the nonlinear elastic system; Q_0 the amplitude of the disturbing factor which may be force or moment; ω - pulsation of the disturbing factor.[6-10]

We change the variable of the form $\tau = pt$, where p is the natural pulsation of the system. In this case we have:

$$\frac{d\tau}{dt} = p, q' = \frac{dq}{d\tau} = \frac{dq}{dt} \cdot \frac{dt}{d\tau} = \dot{q} \frac{1}{p} \quad (3)$$

from where it emerges $\dot{q} = pq'$

For \ddot{q} we perform:

$$q' = \frac{d}{dt} \dot{q} = \frac{d}{dt} (pq') \text{ sau } \ddot{q} = p^2 \frac{d}{d\tau} (q') = p^2 q'' \quad (4)$$

Using the obtained relations to change the variable, the equation becomes:

$$ap^2 q'' + bpq' + k_0(1 + \beta q^2)q = \underline{Q_0} \cos \omega t \quad (5)$$

which we divide by $ap^2 = k_0$ and have:

$$q'' + rq' + (1 + \beta q^2)q = q_0 \cos \omega t \quad (6)$$

where $p = \sqrt{k_0/am}$ is the natural pulsation of the linear elastic system without amortization;

$r = \frac{b}{ap}$ - relative amortization

$q_0 = \frac{Q_0}{k_0}$ - the relative amplitude of the excitation

It is noted $\Omega = \omega/p$ the relative pulsation, so that equation (16.18) can be written as:

$$q'' + rq' + (1 + \beta q^2)q = q_0 \cos \Omega \tau \quad (7)$$

There are analyzed the free unamortized vibrations described by equation (7) where $r = 0$, so that the equation:

$$q'' + q + \beta q^3 = 0 \quad (8)$$

For the nonlinear differential equation (8) choose the solution (in the first approximation) such as:

$$q = A \sin \Omega \tau$$

which inserted together with its derivative in equation (8), leads to

$$-A\Omega^2 \sin \Omega \tau + A \sin \Omega \tau + \beta A^3 \sin^3 \Omega \tau = 0$$

and by transforming the function $\sin^3 \Omega \tau$ and ordering we have

$$A \left(1 - \Omega^2 + \frac{3}{4} \beta A^2 \right) \sin \Omega \tau - \frac{1}{4} \beta A^3 \sin 3\Omega \tau = 0$$

The coefficient of $\sin \Omega \tau$ is canceled and obtained:

$$\Omega^2 = 1 + \frac{3}{4} \beta A^2 \quad (9)$$

The relation (9) represents, in the first approximation, the “amplitude-frequency” characteristic specific to nonlinear systems. It is found that if $\beta > 0$ then the relative pulsation Ω increases by A , and if $\beta < 0$ the relative pulsation decreases by A .

For the second approximation, the solution is chosen

$$q = A \sin \Omega \tau + B \sin 3\Omega \tau \quad (10)$$

which inserted together with its derivative in equation (8) leads to identity.

$$A \left(1 - \Omega^2 + \frac{3}{4} \beta A^2 \right) \sin \Omega \tau + \left(B - 9B\Omega^2 - \frac{1}{4} \beta A^3 \right) \sin 3\Omega \tau + 3A^2 B \beta \sin^2 \Omega \tau \sin 3\Omega \tau + 3AB^2 \beta \sin \Omega \tau \sin^2 3\Omega \tau + \beta B^3 \sin^3 3\Omega \tau = 0$$

in which the coefficient of the function $\sin 3\Omega \tau$ is canceled and it is obtained

$$B = \frac{\beta A^3}{4(1-9\Omega^2)} \quad (11)$$

relation that represents the “amplitude - frequency” characteristic in the second approximation.

If we insert (9) in (11) and approximate $\beta A^2 \cong 0$ then we have

$$B = -\beta \frac{A^3}{32}$$

The analysis of the occurrence of the subharmonic vibrations is based on the fact that the mechanical work provided by the excitation, for the main harmonic in free vibration, must be equal to the mechanical work of dissipation of the system. Thus, for the symmetric nonlinear elastic system, the mechanical work provided by the excitation $Q_0 \cos 3\Omega \tau$ and absorbed by the 3rd order harmonic of the form $q_3 = B \sin 3\Omega \tau$ resonating with the excitation, that is $\Omega = 1$, is of the form [11-15]

$$L_3 = \int_0^{2\pi} Q_0 \cos 3\Omega \tau dq_3 \quad (12)$$

in which it is inserted $dq_3 = 3\Omega B \cos 3\Omega \tau$ and we obtain

$$L_3 = 3BQ_0 \int_0^{2\pi} \cos^2 3\Omega \tau d\tau \quad (13)$$

If $\Omega = 1$, we have

$$L_3 = 3\pi B Q_0 \quad (14)$$

The mechanical work absorbed by the viscous system, on the 1st order component, with the solution $q_1 = A \sin \Omega \tau$, corresponds to the viscous force $F_1 = b \dot{q} = pbq_1'$ and is given by the relation [16-18]

$$L_1 = \int_0^{2\pi} F_1 dq_1 \quad (15)$$

where $F_1 = pbA\Omega \cos \Omega \tau$, and $dq_1 = A\omega \cos \Omega \tau d\tau$
For $\Omega = 1$, the relation (15) becomes

$$L_1 = \int_0^{2\pi} pbA^2 \cos^2 \tau d\tau = \pi pbA^2 \quad (16)$$

the condition of occurrence of the subharmonic of 1/3 order is that $L_3 > L_1$, that is

$$3\pi B Q_0 > \pi pbA^2 \quad (17)$$

in which it is inserted (17) and obtained

$$3Q_0 \left| \frac{\beta A^3}{32} \right| > pbA^2 \quad (18)$$

Taking into account the previously made notations and replacing $b = mpr$ in (18) we obtain

$$r < \frac{3}{32} q_0 A \beta \quad (19)$$

The relation (19) represents the condition of occurrence of the subharmonic vibrations of order 1/3. It is found that subharmonics of order 1/3 cannot be primed (triggered) alone, but only due to excitation to resonance, when amplitude increases a lot. In real physical systems it is difficult to prevent the occurrence of natural vibrations due to external excitations, which is why protection conditions are required for the subharmonic response. This requires that relation (19) be modified as follows: [19-22]

$$r > \frac{3}{32} A_1 q_0 \beta \quad (20)$$

where A_1 is the maximum possible amplitude of the fundamental vibrations in the system.[23-25]

3. CONCLUSIONS

For symmetrical nonlinear elastic systems, in the case of leaning the dynamic equipment, it is highlighted the case of appearance of subharmonic vibrations of order 1/3 when relation $32r < 2q_0 A \beta$ is satisfied.

The avoidance of subharmonic vibrations is possible in this case by adopting the appropriate

hysteretic amortization $r = \frac{b}{ap}$ or $r = \frac{b}{m\omega_n} > \frac{3}{32} A q_0 \beta$

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